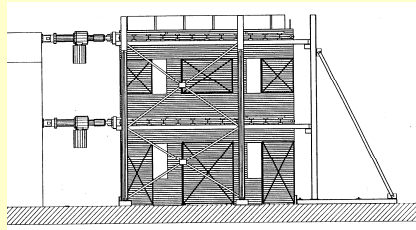
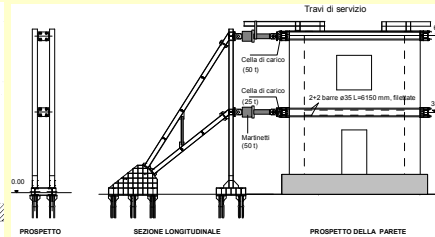


Masonry walls - Modeling in-plane response (to seismic actions)

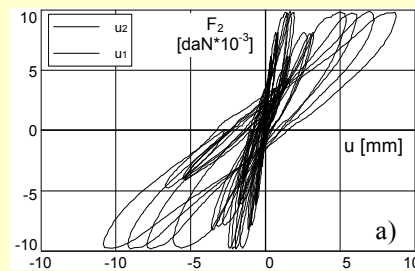
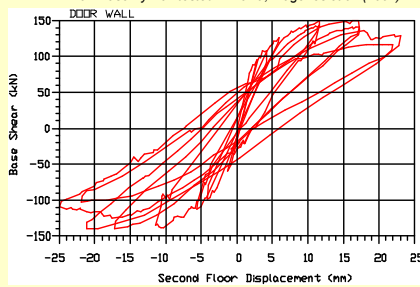
Cyclic horizontal forces, anisotropic damage, hysteretic dissipation, damage localization, inertial vertical forces



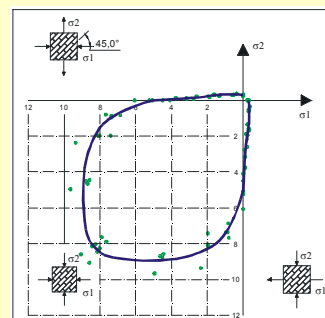
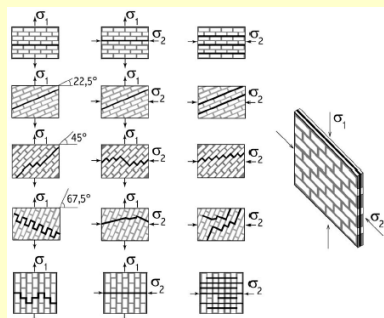
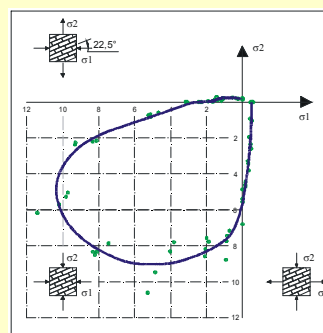
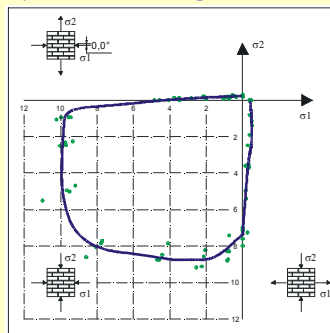
Brick masonry wall tested in Pavia, Magenes et al. (1994).



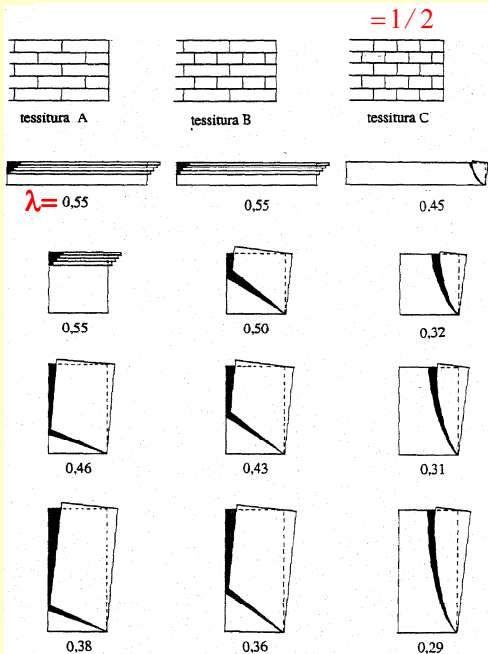
Block masonry wall in S. Sisto (Beolchini et al., 1997).



Anisotropic limit strength domains - Page, 1981



Influence of the brick/block aspect ratio and bond pattern

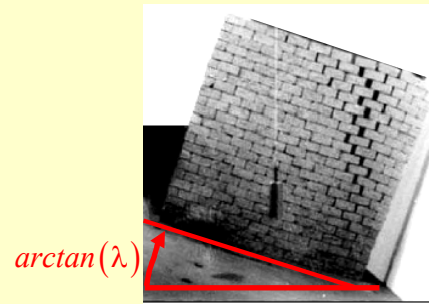


Experimental results
Dry block masonry
Giuffrè *et al.*, 1993

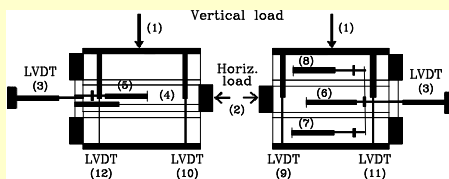
Collapse mechanisms and limit slope angle $\arctan(\lambda)$

for varying:

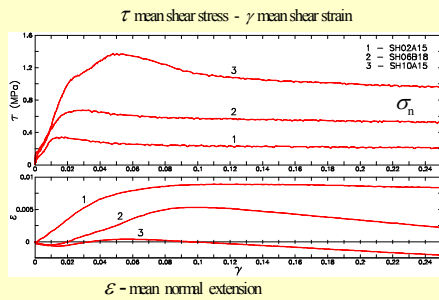
- Block aspect ratio a/b
- Bond pattern
- Wall slenderness



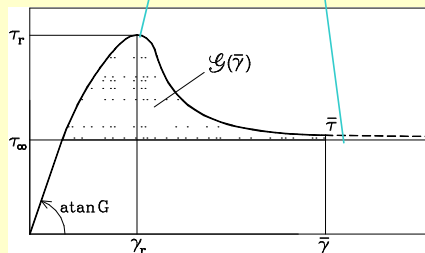
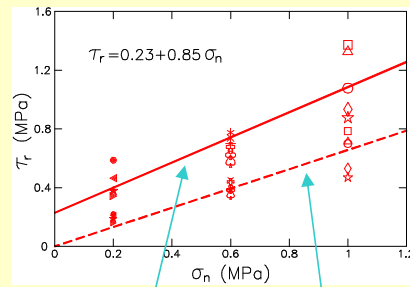
Shear testing on brick-mortar assemblages



Shear test apparatus - Triplet
(Binda *et al.*, 1995).

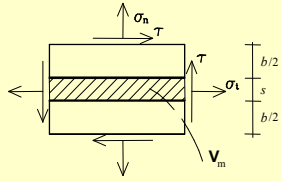


Experimental results



Phenomenological description

Brick-mortar interface model: coupled damage-frictional interface



Macro fields

$$\begin{cases} \boldsymbol{\sigma}_m = \{\sigma_t \ \sigma_n \ \tau\}^t \\ \boldsymbol{\varepsilon}_m^* = \{0 \ \varepsilon_m^* \ \gamma_m^*\}^t \quad \text{Inelastic strain} \\ \boldsymbol{\varepsilon}_m = \{0 \ \varepsilon_m \ \gamma_m\}^t \quad \text{Total strain} \end{cases}$$

$$\begin{cases} \varepsilon_m^* = h(\alpha_m) H(\sigma_n) \sigma_n \\ \gamma_m^* = k(\alpha_m) (\tau - f) \end{cases}$$

$$\boldsymbol{\varepsilon}_m = \mathbf{K}_m \boldsymbol{\sigma}_m + \boldsymbol{\varepsilon}_m^*$$

α_m damage variable
f friction

Conjugate variables

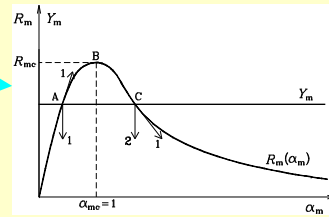
$$Y_m = \frac{1}{2} h'(\alpha_m) H(\sigma_n) \sigma_n^2 + \frac{1}{2} k'(\alpha_m) (\tau - f)^2, \quad \dot{\gamma}_m^*$$

Damage evolution

$$\begin{cases} \phi_{dm} = Y_m - R_m \leq 0 \\ \phi_{dm} = 0, \dot{\phi}_{dm} \leq 0, \dot{\alpha}_m \geq 0, \dot{\phi}_{dm} \dot{\alpha}_m = 0 \end{cases}$$

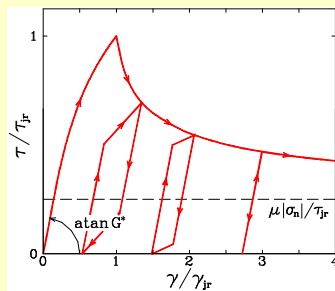
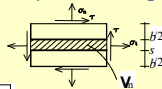
Sliding

$$\begin{cases} \phi_s = |f| + \mu \sigma_n \leq 0 \\ \dot{\gamma}_m^* = v \dot{\lambda}, \quad \dot{\lambda} \geq 0 \quad v = \frac{f}{|f|} \end{cases}$$

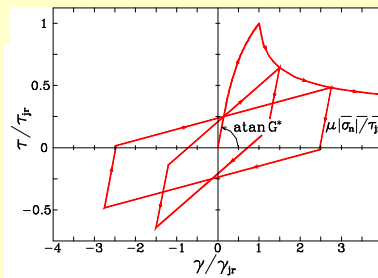


Gambarotta e Lagomarsino, 1997

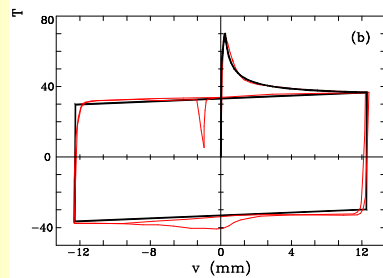
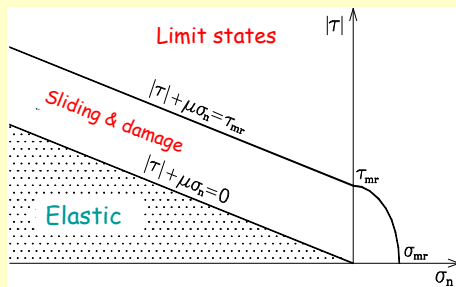
Brick-mortar interface model: coupled damage-frictional interface

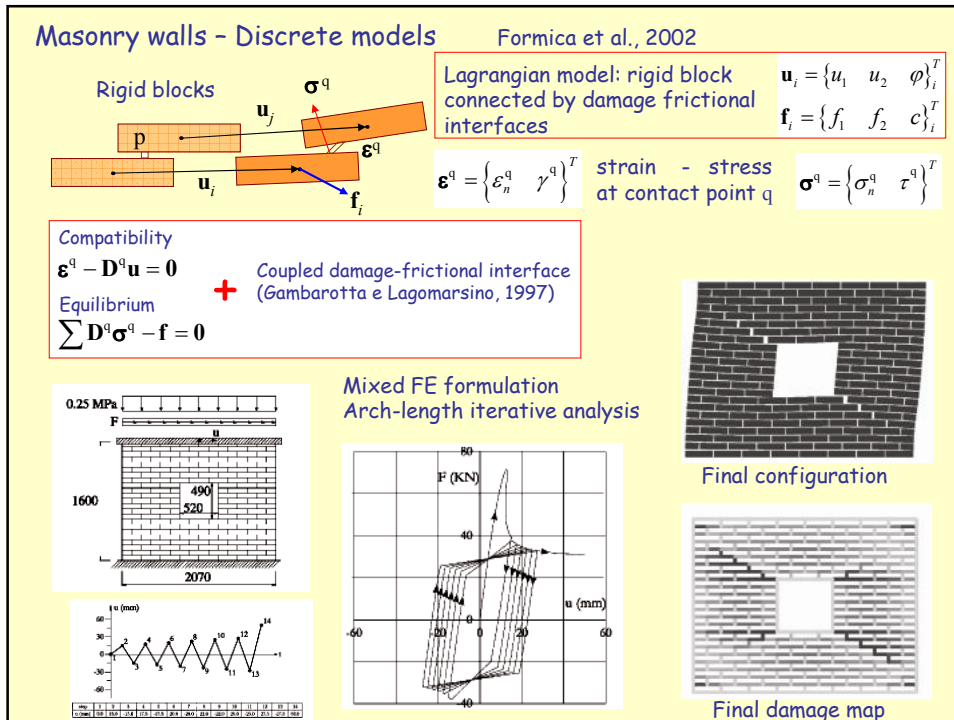
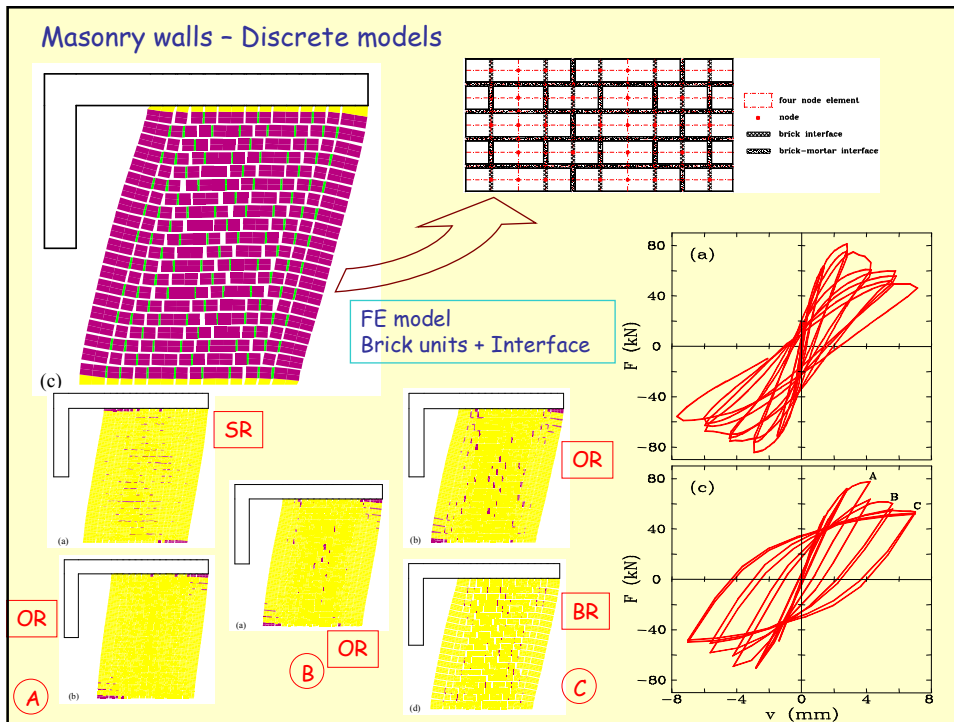


Hysteretic damage



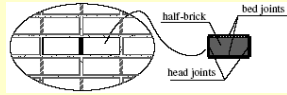
Simulation of experimental results (Atkinson et al)



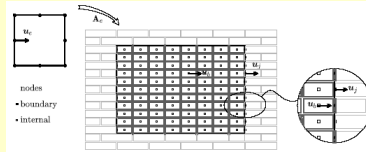


Discrete models - Multigrid/Multilevel approach

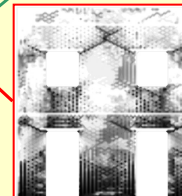
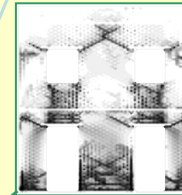
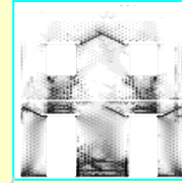
Casciaro *et al*, CMAME, 2007



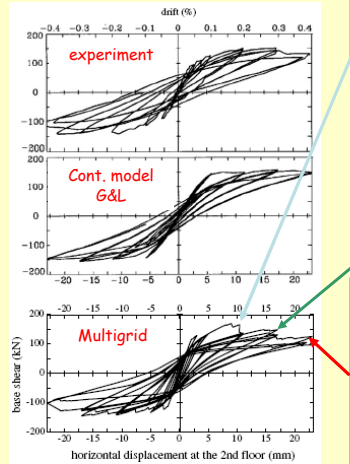
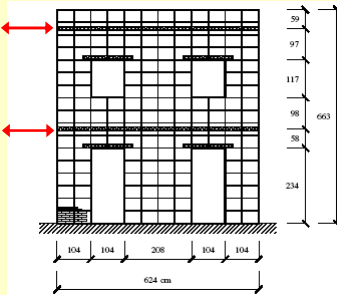
Rigid blocks + Interfaces



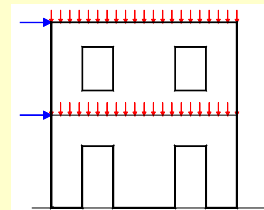
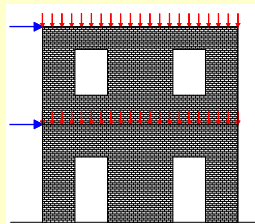
Damage



Simulation of Pavia experiment



Large masonry shear walls - 2D Cauchy equivalent continuum



Meso fields $\sigma, u, \varepsilon, \zeta$

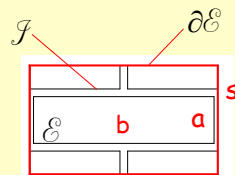
$$\mathbf{u}(\mathbf{x}) = \mathbf{E}\mathbf{x} + \mathbf{u}_{\text{per}}$$

\mathbf{u}_{per} periodic on $\partial\mathcal{E}$

$$\text{div}\boldsymbol{\sigma} = \mathbf{0} \text{ in } \mathcal{E}$$

$\boldsymbol{\sigma}\mathbf{n}$ antiperiodic on $\partial\mathcal{E}$

$$\|\boldsymbol{\sigma}\|_{\mathbf{n}} = \mathbf{0} \text{ su } \bar{\mathcal{J}}$$



Periodic RVE

Macro fields $\Sigma, \mathbf{E}, \mathbf{Z}$

$$\boldsymbol{\Sigma} = \frac{1}{A} \int_{\partial\mathcal{E}} \mathbf{x} \otimes \mathbf{t} ds$$

$$\mathbf{E} = \frac{1}{A} \int_{\partial\mathcal{E}} \text{sym}(\mathbf{u} \otimes \mathbf{n}) ds$$

Meso - constitutive equations

Brick units $\boldsymbol{\sigma}_b \leftrightarrow \boldsymbol{\varepsilon}_b, \zeta_b$

Mortar $\boldsymbol{\sigma}_m \leftrightarrow \boldsymbol{\varepsilon}_m, \zeta_m$

Interface $\boldsymbol{\sigma}_i \leftrightarrow \boldsymbol{\varepsilon}_i, \zeta_i$

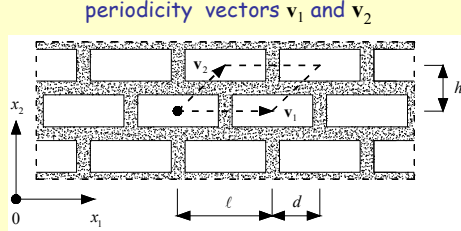
ζ internal variables

Macro - constitutive equations

$\boldsymbol{\Sigma} \leftrightarrow \mathbf{E}, \mathbf{Z}$

\mathbf{Z} internal variables

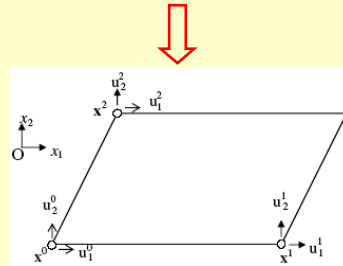
Homogenization of periodic masonry Anthoine, 1995, 2007



strain-periodic cinematically admissible displacement field

$$\mathbf{u}(\mathbf{x}) = \tilde{\mathbf{E}}\mathbf{x} + \tilde{\mathbf{W}}\mathbf{x} + \mathbf{u}^p(\mathbf{x})$$

$$\tilde{\mathbf{W}} = \mathbf{0}, \mathbf{u}^0 = \mathbf{0}$$



$$\begin{aligned} \tilde{\mathbf{T}} &= \frac{1}{A} \int \mathbf{t} \otimes \mathbf{x} \, ds = \\ &= \frac{1}{A} \sum_{i=0}^2 \mathbf{f}_i \otimes \mathbf{x}_i = \mathbf{C}_{\text{hom}} \tilde{\mathbf{E}} \end{aligned}$$

$$\begin{aligned} u_1(\mathbf{x} + \mathbf{v}_1) - u_1(\mathbf{x}) &= u_1^1, \\ u_2(\mathbf{x} + \mathbf{v}_1) - u_2(\mathbf{x}) &= u_2^1, \\ u_1(\mathbf{x} + \mathbf{v}_2) - u_1(\mathbf{x}) &= \frac{d}{\ell} u_1^1 + \frac{h}{\ell} u_2^1 = u_1^2, \\ u_2(\mathbf{x} + \mathbf{v}_2) - u_2(\mathbf{x}) &= u_2^2, \end{aligned}$$

2D micropolar orthotropic continuum models - Periodic rigid block masonry Besdo, Sulem & Mühlhaus, Di Carlo, Rizzi, Trovalusci & Masiani, Sulem & Vardoulakis

-> Trovalusci & Masiani, 2003, 05 Running bond 2D

- Lagrangian model: dofs block A $\mathbf{u}_A, \mathbf{W}_A$

- Homogeneity of the generalized displacements gradients $\mathbf{H} - \mathbf{K}$ in the RVE
 $\mathbf{H} = \text{grad} \mathbf{u}, \mathbf{K} = \text{grad} \mathbf{W}, \mathbf{\Gamma} = \mathbf{H} - \mathbf{W}$

- Relative generalized displacement at interface C

$$\begin{cases} \mathbf{u}_C = \mathbf{\Gamma} \mathbf{v}_{BA} + (\mathbf{K} \mathbf{v}_{BX}) \mathbf{v}_{CB} - (\mathbf{K} \mathbf{v}_{AX}) \mathbf{v}_{CA} \\ \mathbf{W}_C = \mathbf{K} \mathbf{v}_{BA} \end{cases}$$

- $\mathbf{t}_C, \mathbf{T}_C$ elastic generalized contact forces at interface C $\mathbf{u}_A, \mathbf{W}_A$

- Equivalence of the work expended per unit area in the Lagrangian and in the Polar continuum

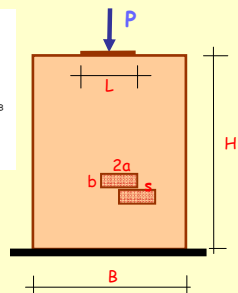
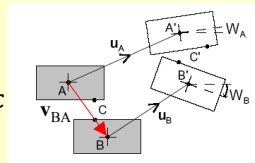
$$\pi_{LAG}(\mathbf{\Gamma}, \mathbf{K}) = \frac{1}{A} \sum_C \left(\mathbf{t}_C \cdot \mathbf{u}_C + \frac{1}{2} \mathbf{T}_C \cdot \mathbf{W}_C \right) = \pi_{HOM}(\mathbf{\Gamma}, \mathbf{K}) = \mathbf{\Sigma} \cdot \mathbf{\Gamma} + \frac{1}{2} \mathbf{M} \cdot \mathbf{K}$$

- Average stress $\mathbf{\Sigma}$ and couple stress \mathbf{M} tensors

$$\begin{cases} \mathbf{\Sigma} = \mathbf{Y} \mathbf{\Gamma} \\ \mathbf{M} = \mathbf{C} \mathbf{K} \end{cases}$$

- Constitutive equation in the case of centro-symmetric unit cell =

$$\begin{aligned} Y_{1111} &= \delta^{-1} (k_t^+ + 2k_n^+), & \delta &= b/a \\ Y_{2222} &= \delta k_n^+ \\ Y_{1212} &= \delta k_t^+ \\ Y_{2121} &= \delta^{-1} (k_n^+ + 2k_t^+) \\ C_{31} &= a^2 (4k_n^+ + 3k_t^+ \delta^2 + 2k_n^+ \delta^2 + 12k_t^+) / 12\delta \\ C_{32} &= a^2 \delta^2 (4k_n^+ + 3k_t^+ \delta^2 + 12k_t^+) / 12 \end{aligned}$$

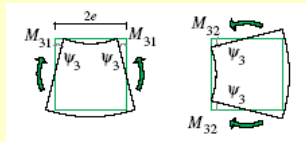


2D micropolar continuum models - Periodic masonry with elastic units and mortar

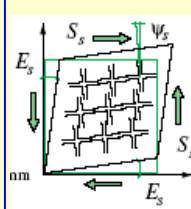
-> Casolo, 2004,06 - Running bond

$$\begin{aligned} \Gamma_{11} &= E_{11} & E_{ii} &= U_{i,i} \\ \Gamma_{22} &= E_{22} & E_{12} &= \frac{1}{2}(U_{1,2} + U_{2,1}) \\ \Gamma_{12} &= E_{12} + \theta & \theta &= \phi - \Omega \\ \Gamma_{21} &= E_{12} - \theta & \Omega &= \frac{1}{2}(U_{2,1} - U_{1,2}) \end{aligned}$$

$$\begin{Bmatrix} \Sigma_{11} \\ \Sigma_{22} \\ \Sigma_{12} \\ \Sigma_{21} \\ M_1 \\ M_2 \end{Bmatrix} = \begin{bmatrix} Y_{1111} & Y_{1122} & 0 & 0 & 0 & 0 \\ Y_{2211} & Y_{2222} & 0 & 0 & 0 & 0 \\ 0 & 0 & Y_{1212} & Y_{1221} & 0 & 0 \\ 0 & 0 & Y_{2112} & Y_{2121} & 0 & 0 \\ 0 & 0 & 0 & 0 & C_{3131} & 0 \\ 0 & 0 & 0 & 0 & 0 & C_{3232} \end{bmatrix} \begin{Bmatrix} E_{11} \\ E_{22} \\ E_{12} + \theta \\ E_{21} - \theta \\ K_1 \\ K_2 \end{Bmatrix}$$



(a)



Prescribed displacements at the boundary of the RVE

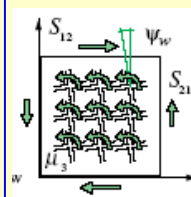
$$\mathbf{u} = E_{12} \text{sym}(\mathbf{e}_1 \otimes \mathbf{e}_2) \mathbf{x} + \mathbf{u}^p$$

Evaluate:

Mean (symmetric) stress Σ_{12}^s

Block rotation (specifically defined) $\phi_s = \Psi_s$

(b)

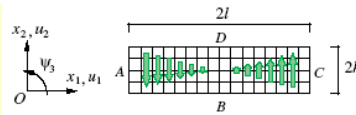


Prescribed rotation $\phi_w = \Psi_w$ of each block with periodic boundary condition $\mathbf{u} = \mathbf{u}^p$

Evaluate:

Mean couple density

$$\mu_3 = \Sigma_{21} - \Sigma_{12} \quad (E_{12} = \Omega = 0)$$

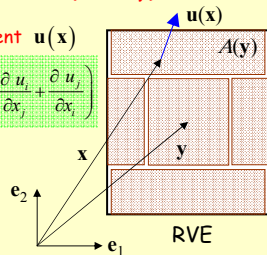


2D micropolar homogenization - Forest & Sab, 1998

Heterogeneous (Cauchy) medium

displacement $\mathbf{u}(\mathbf{x})$

$$\varepsilon_{ij}(\mathbf{x}) = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$



Generalized (Cosserat) continuum

displacement $\mathbf{v}(\mathbf{y}), \phi(\mathbf{y})$

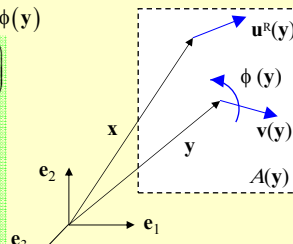
$$E_{ij}(\mathbf{y}) = \frac{1}{2} \left(\frac{\partial v_i}{\partial y_j} + \frac{\partial v_j}{\partial y_i} \right)$$

$$\Omega(\mathbf{y}) = \frac{1}{2} \left(\frac{\partial v_2}{\partial y_1} - \frac{\partial v_1}{\partial y_2} \right)$$

$$\theta(\mathbf{y}) = \phi(\mathbf{y}) - \Omega(\mathbf{y})$$

$$K_i(\mathbf{y}) = \frac{\partial \phi}{\partial y_i}$$

$$\mathbf{u}^R(\mathbf{x}) = \mathbf{v}(\mathbf{y}) + \phi(\mathbf{y}) \mathbf{e}_3 \times (\mathbf{x} - \mathbf{y}) \quad \text{local rigid motion}$$



Characterization of the rigid motion that best fits the actual displacement field

$$\min \int_A |\mathbf{u}(\mathbf{x}) - \mathbf{v}(\mathbf{y}) + \phi(\mathbf{y}) \mathbf{e}_3 \times (\mathbf{x} - \mathbf{y})|^2 da$$

$$\mathbf{v}(\mathbf{y}) = \frac{1}{A} \int_A \mathbf{u} da = \langle \mathbf{u} \rangle_A \quad J_p = \int_A |\mathbf{x} - \mathbf{y}|^2 da$$

$$\phi(\mathbf{y}) = \frac{1}{J_p} \int_A [u_2(x_1 - y_1) - u_1(x_2 - y_2)] da$$

Generalized displacement components

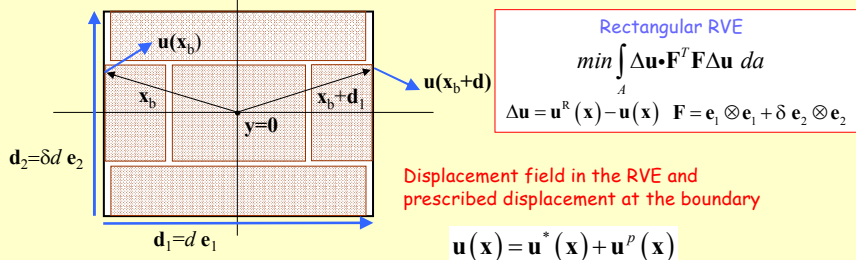
Macro-displ gradient

$$\frac{\partial v_i}{\partial y_j} = \frac{1}{A} \int_A u_i n_j da = \left\langle \frac{\partial u_i}{\partial x_j} \right\rangle_A$$

Curvature

$$K_i = \frac{\partial \phi}{\partial y_i} = \frac{A}{J_p} \left[\frac{1}{A} \int_A \left(\frac{\partial u_2}{\partial x_j} x_j - \frac{\partial u_1}{\partial x_j} x_j \right) da - \frac{\partial v_2}{\partial y_j} y_j + \frac{\partial v_1}{\partial y_j} y_j \right]$$

Homogenization criterion based on equivalence of strain energy density in RVE



- $\mathbf{u}^p(\mathbf{x})$ displacement field fulfilling the periodicity conditions at the boundary C

$$\mathbf{u}^p(\mathbf{x}_b) = \mathbf{u}^p(\mathbf{x}_b + \mathbf{d}_i) \quad i=1,2 \quad \mathbf{x}_b \in C$$

- $\mathbf{u}^*(\mathbf{x})$ displacement polynomial of grade 3 satisfying the following conditions:

$\mathbf{E}^*(\mathbf{y}) = \boldsymbol{\varepsilon}^*(\mathbf{y})$ of grade 2 - **scale invariant**

$\phi^*(\mathbf{y})$ affine function

$$\mathbf{u}^*(\mathbf{x}) = \begin{cases} u_1^* = B_{11}\tilde{x}_1 + B_{12}\tilde{x}_2 - C_{23}\tilde{x}_2^2 + 2C_{13}\tilde{x}_1\tilde{x}_2 + D_{12}(\tilde{x}_2^3 - 3\tilde{x}_1^2\tilde{x}_2) \\ u_2^* = B_{21}\tilde{x}_1 + B_{22}\tilde{x}_2 - C_{13}\tilde{x}_1^2 + 2C_{23}\tilde{x}_1\tilde{x}_2 - D_{12}(\tilde{x}_1^3 - 3\tilde{x}_1\tilde{x}_2^2) \end{cases} \quad \tilde{x}_i = x_i/d$$

Boundary conditions on the RVE $\mathbf{u}(\mathbf{x}_b + \mathbf{d}_i) - \mathbf{u}(\mathbf{x}_b) = \mathbf{u}^*(\mathbf{x}_b + \mathbf{d}_i) - \mathbf{u}^*(\mathbf{x}_b) \quad i=1,2 \quad \mathbf{x}_b \in C$

Cauchy continuum homogenization $\mathbf{u}^*(\mathbf{x}) = \mathbf{E}^* \mathbf{x}$ \mathbf{E}^* mean strain tensor

Components of generalized strain at the centre $\mathbf{y}=\mathbf{0}$ of the RVE

- 3 components of the symm part of the strain tensor $E_{ij}(\mathbf{y}=\mathbf{0}) = E_{ij}^*(\mathbf{y}=\mathbf{0}) = \bar{E}_{ij}$
- 1 comp. of the skew part of the strain tensor $\theta(\mathbf{y}=\mathbf{0}) = \phi(\mathbf{y}=\mathbf{0}) - \Omega(\mathbf{y}=\mathbf{0}) = \bar{\theta}$
- 2 components of the torsion curvature tensor $\mathbf{K}(\mathbf{y}=\mathbf{0}) = \{\bar{K}_1 \quad \bar{K}_2\}^T = \bar{\mathbf{K}}$

Assuming $\Omega(\mathbf{y}=\mathbf{0})=0 \Rightarrow B_{12}=B_{21}$

$$B_{11}/d = \bar{E}_{11}, \quad B_{22}/d = \bar{E}_{22}, \quad B_{12}/d = \bar{E}_{12}$$

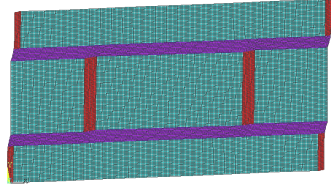
$$\frac{\delta^2}{10d} D_{12} + \frac{6}{\delta^3 d^4} \int_A (\delta^2 u_2^p x_1 - u_1^p x_2) da = \bar{\theta}$$

$$-\frac{2}{d^2} C_{13} + \frac{6}{\delta d^4} \int_A u_{2,1}^p x_1 da = \bar{K}_1$$

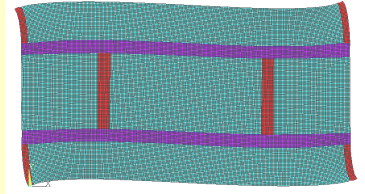
$$\frac{2}{d^2} C_{23} - \frac{6}{\delta^3 d^4} \int_A u_{1,2}^p x_2 da = \bar{K}_2$$

Homogenization of a regular pattern - English Bond

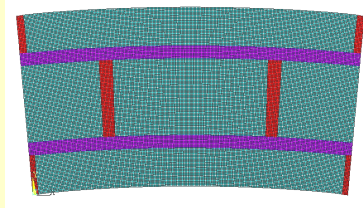
$B_{12} = 1 \Rightarrow E_{12} \neq 0, \theta = \theta^p = \phi \neq 0, \Omega = 0$



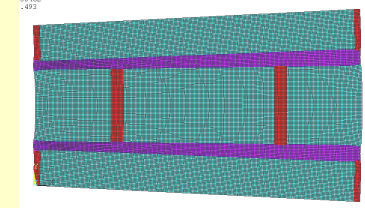
$D_{12} = 1 \Rightarrow E_{12} = \Omega = 0, \theta \neq 0$



$C_{31} = 1 \Rightarrow K_1 \neq 0, E_{ij} = \theta = 0$

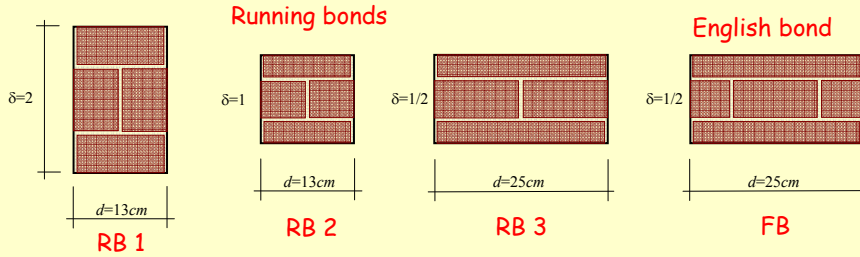


$C_{32} = 1 \Rightarrow K_2 \neq 0, E_{ij} = \theta = 0$



$E_b = 5000MPa, E_m = 500MPa, \nu_b = \nu_m = 0.1, s = 10mm$

Elastic generalized moduli for varying masonry patterns



$E_b = 5000MPa, E_m = 500MPa, \nu_b = \nu_m = 0.1, s = 10mm$

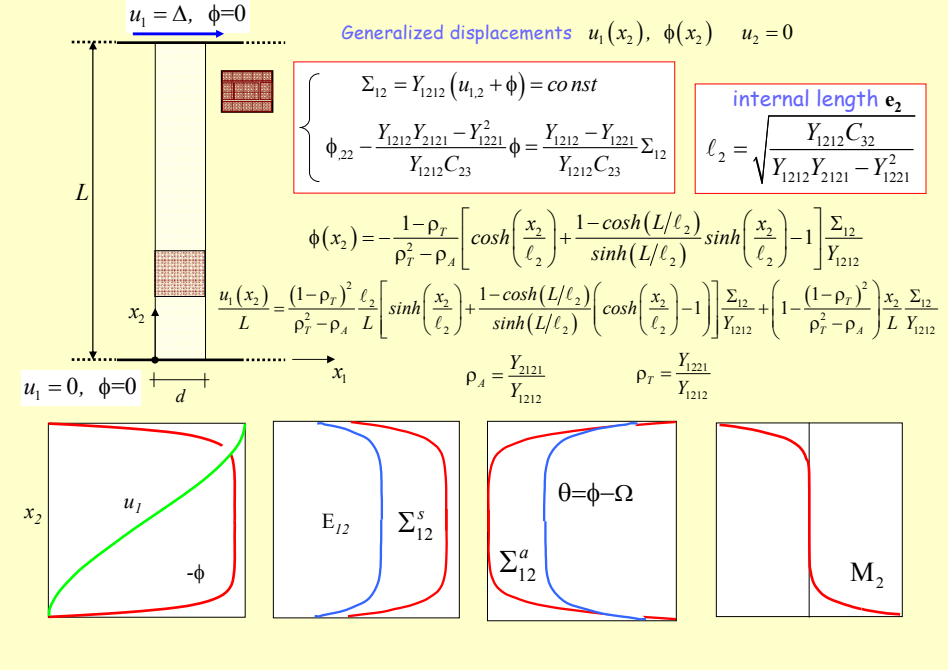
	Y_{1111} (Mpa)	Y_{2222} (Mpa)	Y_{1122} (Mpa)	Y_{1212} (Mpa)	Y_{2121} (Mpa)	Y_{1221} (Mpa)	C_{31} (N)	C_{32} (N)	ℓ_1 (mm)	ℓ_2 (mm)
RB1	2,865E+03	2,814E+03	2,205E+02	1,407E+04	1,021E+04	-9,899E+03	1,529E+07	4,169E+06	58	36
RB2	2,776E+03	2,098E+03	1,845E+02	7,130E+03	5,141E+03	-4,351E+03	4,553E+06	3,080E+06	36	35
RB3	3,466E+03	2,154E+03	1,990E+02	6,741E+03	4,770E+03	-3,838E+03	5,937E+06	1,203E+07	40	68
FB	3,101E+03	2,126E+03	1,911E+02	7,143E+03	4,461E+03	-3,876E+03	5,816E+06	1,180E+07	39	71

Y_{ijhk} scale invariant
 $\frac{C_{3i}^d}{C_{3i}^D} = \left(\frac{d}{D}\right)^2$

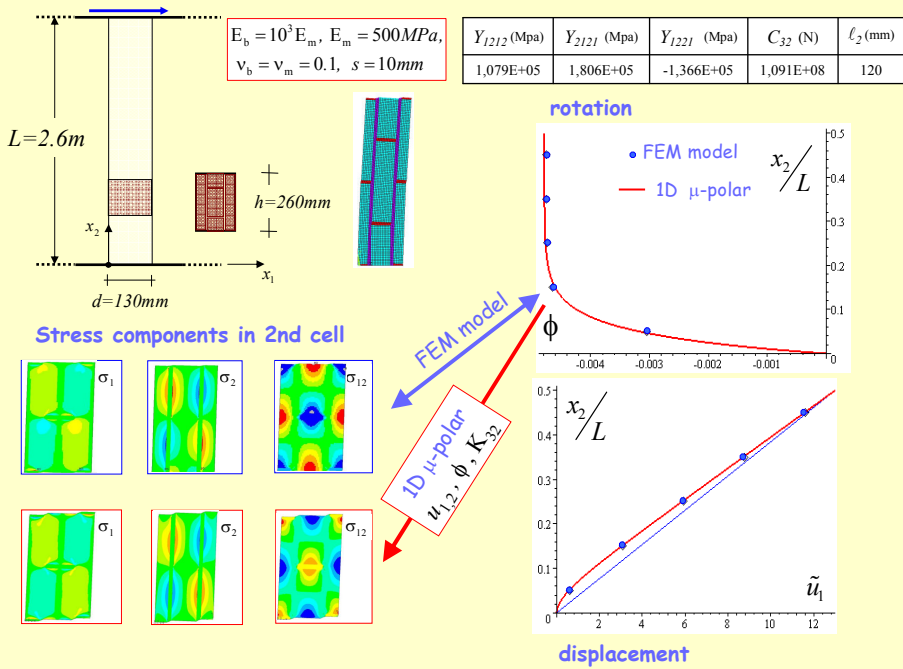
internal lengths
 $\ell_1 = \sqrt{\frac{Y_{2121}C_{31}}{Y_{1212}Y_{2121} - Y_{1221}^2}} \quad \ell_2 = \sqrt{\frac{Y_{1212}C_{32}}{Y_{1212}Y_{2121} - Y_{1221}^2}}$

Bacigalupo, Gambarotta, 2008

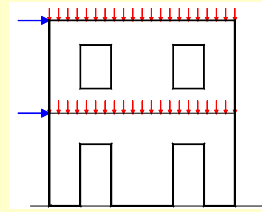
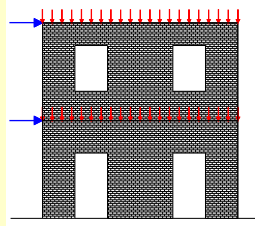
Shearing of a masonry walls - Boundary shear layer model



Evaluation of the omogenization of running bond masonry - vertical cell



Damage in masonry shear walls



Meso fields $\sigma, u, \varepsilon, \zeta$

$$\mathbf{u}(\mathbf{x}) = \mathbf{E}\mathbf{x} + \mathbf{u}_{\text{per}}$$

\mathbf{u}_{per} periodic on $\partial\mathcal{E}$

$$\text{div}\boldsymbol{\sigma} = \mathbf{0} \text{ in } \mathcal{E}$$

$\boldsymbol{\sigma}\mathbf{n}$ antiperiodic on $\partial\mathcal{E}$

$$\|\boldsymbol{\sigma}\|_{\mathbf{n}} = \mathbf{0} \text{ su } \bar{\mathcal{J}}$$

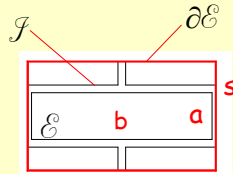
Meso - constitutive equations

Brick units $\sigma_b \leftrightarrow \varepsilon_b, \zeta_b$

Mortar $\sigma_m \leftrightarrow \varepsilon_m, \zeta_m$

Interface $\sigma_i \leftrightarrow \varepsilon_i, \zeta_i$

ζ internal variables



Periodic RVE

Macro fields $\Sigma, \mathbf{E}, \mathbf{Z}$

$$\boldsymbol{\Sigma} = \frac{1}{A} \int_{\partial\mathcal{E}} \mathbf{x} \otimes \mathbf{t} ds$$

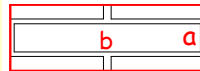
$$\mathbf{E} = \frac{1}{A} \int_{\partial\mathcal{E}} \text{sym}(\mathbf{u} \otimes \mathbf{n}) ds$$

Macro - constitutive equations

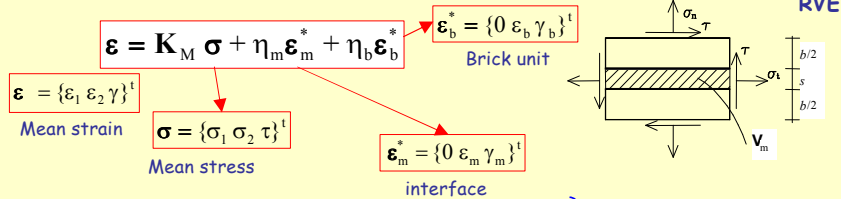
$$\boldsymbol{\Sigma} \leftrightarrow \mathbf{E}, \mathbf{Z}$$

\mathbf{Z} internal variables

Continuum damage-frictional model - Layered model - $b \gg a$



(Gambarotta & Lagomarsino, 1997,
Uva & Salerno, 2006)



Interface - damage & friction

$$\varepsilon_m = c_{mn} \alpha_m H(\sigma_2) \sigma_2$$

$$\gamma_m = c_{mt} \alpha_m (\tau - f)$$

Internal variables: α_m damage & f interface friction

$$\text{Conjugate variables} \implies Y_m = \frac{1}{2} c_{mn} H(\sigma_2) \sigma_2^2 + \frac{1}{2} c_{mt} (\tau - f)^2, \gamma_m$$

Brick unit - damage in compression

$$\varepsilon_b = c_{bn} \alpha_b H(-\sigma_2) \sigma_2$$

$$\gamma_b = c_{bt} \alpha_b \tau$$

Internal variable: α_b brick damage

$$\text{Conjugate variable} \implies Y_b = \frac{1}{2} c_{bn} H(-\sigma_2) \sigma_2^2 + \frac{1}{2} c_{bt} \tau^2$$

Limit conditions:

• Damage

$$\phi_{dm} = Y_m - R_m(\alpha_m) \leq 0$$

$$\phi_{db} = Y_b - R_b(\alpha_b) \leq 0$$

• Friction

$$\phi_s = |f| + \mu \sigma_2 \leq 0$$

sliding

$$\dot{\gamma}_m = v \dot{\lambda}, \dot{\lambda} \geq 0$$

$$v = f/|f|$$

4. Continuum damage-friction model

Layered micro-model

(Gambarotta e Lagomarsino, 1997)

Evolution of the internal variables

$\sigma_2 \geq 0$

Opened interface

$$\phi_{dm} = \frac{1}{2} c_{mn} \sigma_2^2 + \frac{1}{2} c_{mt} \tau^2 - R_m(\alpha_m) \leq 0$$

$$\phi_{db} = \frac{1}{2} c_{bt} \tau^2 - R_b(\alpha_b) \leq 0$$

$$\begin{Bmatrix} \dot{\phi}_{dm} \\ \dot{\phi}_{db} \end{Bmatrix} = \begin{bmatrix} R'_m & 0 \\ 0 & R'_b \end{bmatrix} \begin{Bmatrix} \dot{\alpha}_m \\ \dot{\alpha}_b \end{Bmatrix} + \begin{Bmatrix} c_{mn} \sigma_2 \dot{\sigma}_2 + c_{mt} \tau \dot{\tau} \\ c_{bt} \tau \dot{\tau} \end{Bmatrix} \leq 0$$

$$\begin{Bmatrix} \dot{\phi}_{dm} & \dot{\phi}_{db} \end{Bmatrix} \begin{Bmatrix} \dot{\alpha}_m & \dot{\alpha}_b \end{Bmatrix}^t = 0 \quad \begin{Bmatrix} \dot{\alpha}_m & \dot{\alpha}_b \end{Bmatrix}^t \geq 0$$

$\sigma_2 < 0$

Closed interface

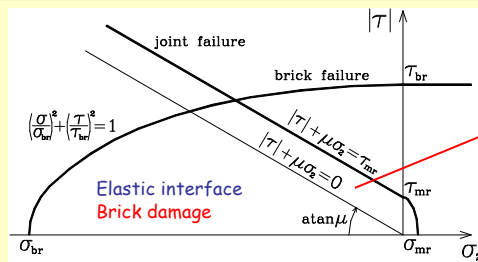
$$\phi_{dm} = \frac{1}{2} \frac{\gamma_m^2}{c_{mt} \alpha_m^2} - R_m(\alpha_m) \leq 0$$

$$\phi_s = \left| \tau - \frac{\gamma_m}{c_{mt} \alpha_m} \right| + \mu \sigma_2 \leq 0$$

$$\phi_{db} = \frac{1}{2} c_{bn} \sigma_2^2 + \frac{1}{2} c_{bt} \tau^2 - R_b(\alpha_b) \leq 0$$

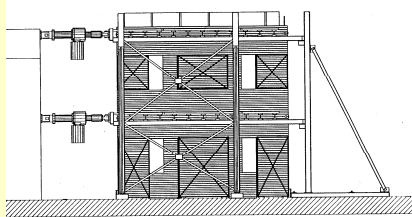
$$\begin{Bmatrix} \dot{\phi}_{dm} \\ \dot{\phi}_s \\ \dot{\phi}_{db} \end{Bmatrix} = \begin{bmatrix} -\frac{\gamma_m^2}{c_{mt} \alpha_m^3} - R'_m & \frac{\gamma \dot{\gamma}_m}{c_{mt} \alpha_m^2} & 0 \\ \frac{\gamma \dot{\gamma}_m}{c_{mt} \alpha_m^2} & -1 & 0 \\ 0 & 0 & R'_b \end{bmatrix} \begin{Bmatrix} \dot{\alpha}_m \\ \dot{\lambda} \\ \dot{\alpha}_b \end{Bmatrix} + \begin{Bmatrix} 0 \\ v \tau + \mu \dot{\sigma}_2 \\ c_{bn} \sigma_2 \dot{\sigma}_2 + c_{bt} \tau \dot{\tau} \end{Bmatrix} \leq 0$$

$$\begin{Bmatrix} \dot{\phi}_{dm} & \dot{\phi}_s & \dot{\phi}_{db} \end{Bmatrix} \begin{Bmatrix} \dot{\alpha}_m & \dot{\lambda} & \dot{\alpha}_b \end{Bmatrix}^t = 0 \quad \begin{Bmatrix} \dot{\alpha}_m & \dot{\lambda} & \dot{\alpha}_b \end{Bmatrix}^t \geq 0$$



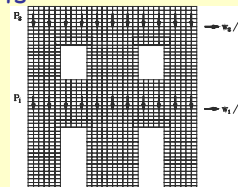
Damage in the interface and brick units

Large shear walls - simulation of experimental results

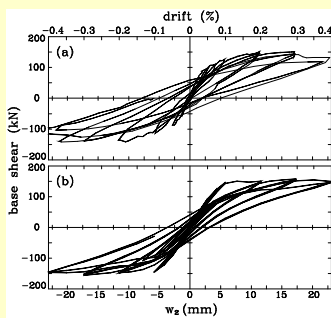
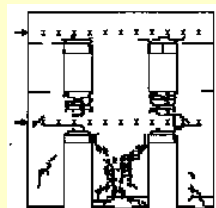


Brick masonry wall tested in Pavia, Magenes et al. (1994).

FEM model



Damage (exp)



Cyclic response of the door walk a) experimental; b) numerical simulation.

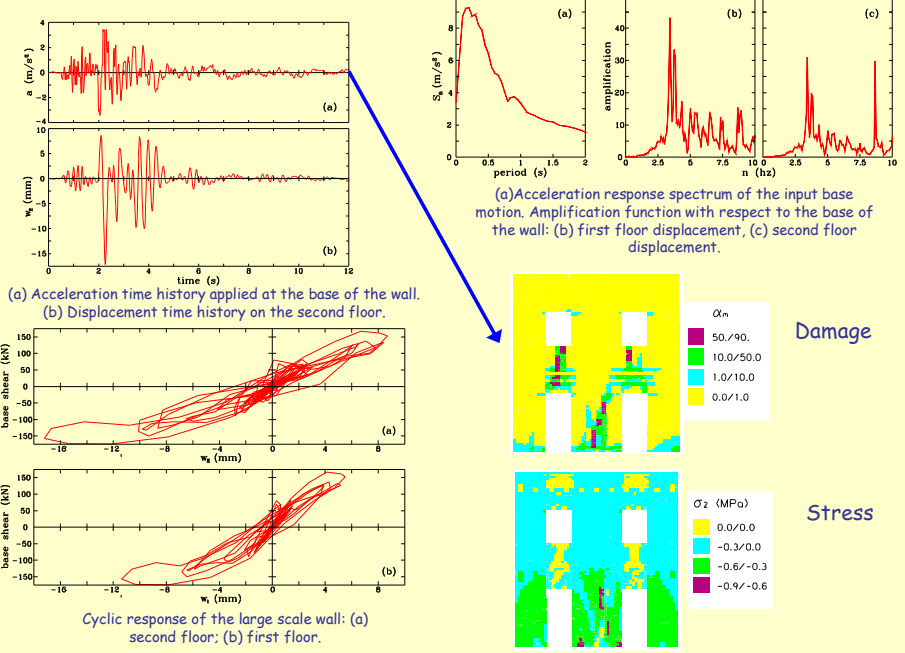
simul



Damage (simul)

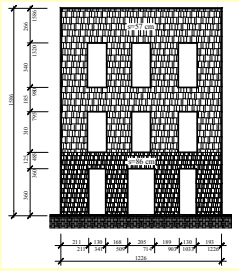


Large shear walls - dynamic response to ground motion

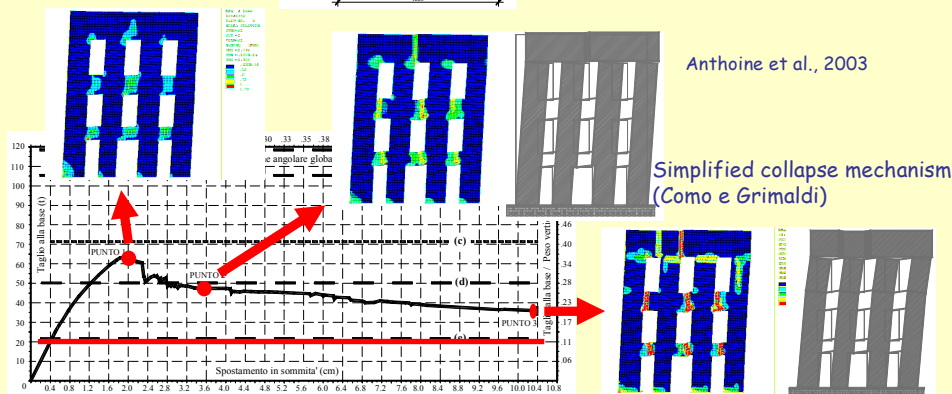
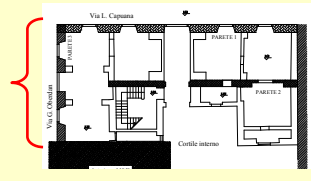


Large shear walls - response to horizontal forces

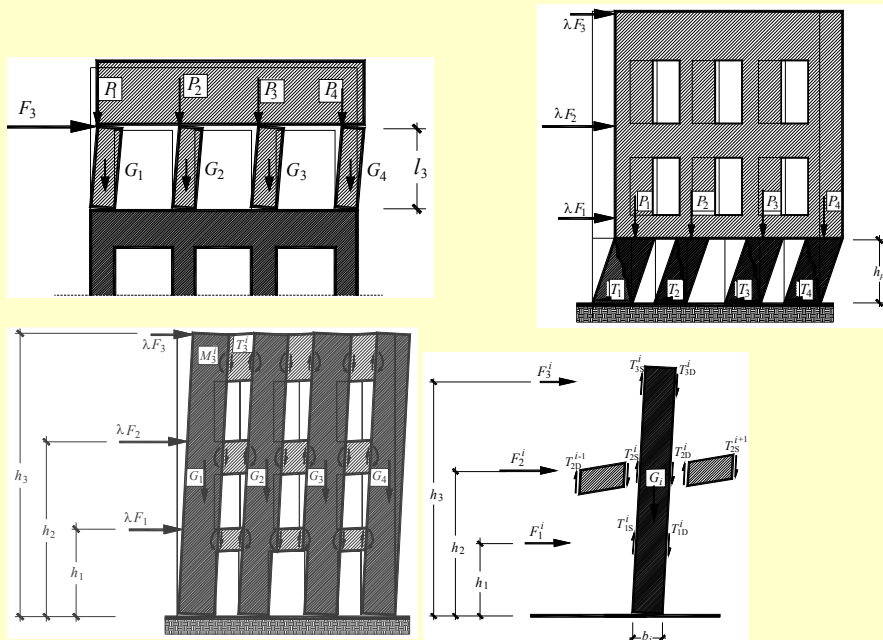
Horizontal forces superimposed on vertical dead loads



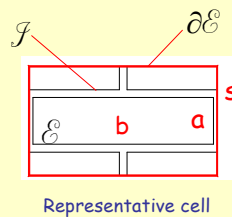
Brencich et al, 2001
 Masonry building in Catania
 GNDT



4. Large shear walls - simplified approaches



Modelling inelastic response of shear walls - Cauchy continuum models

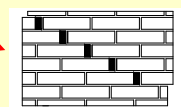
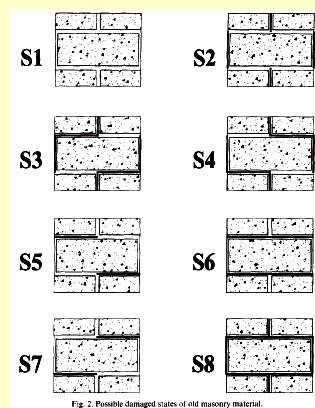


Limit Analysis

Alpa & Monetto, 1994
De Felice & de Buhan, 1995
Milani et al, 2006

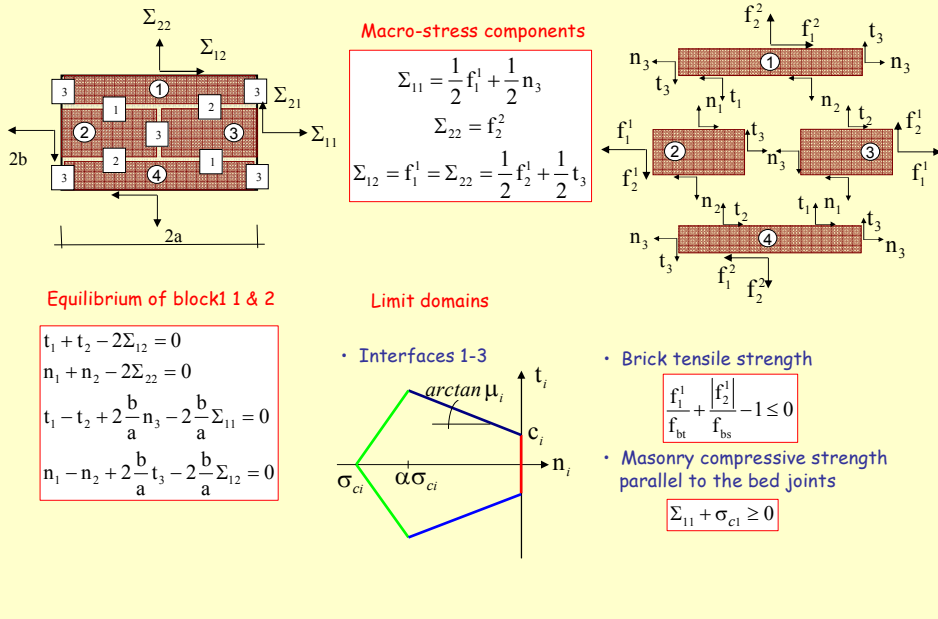
Damage models

Luciano e Sacco, 1997
Massart et al, 2004-

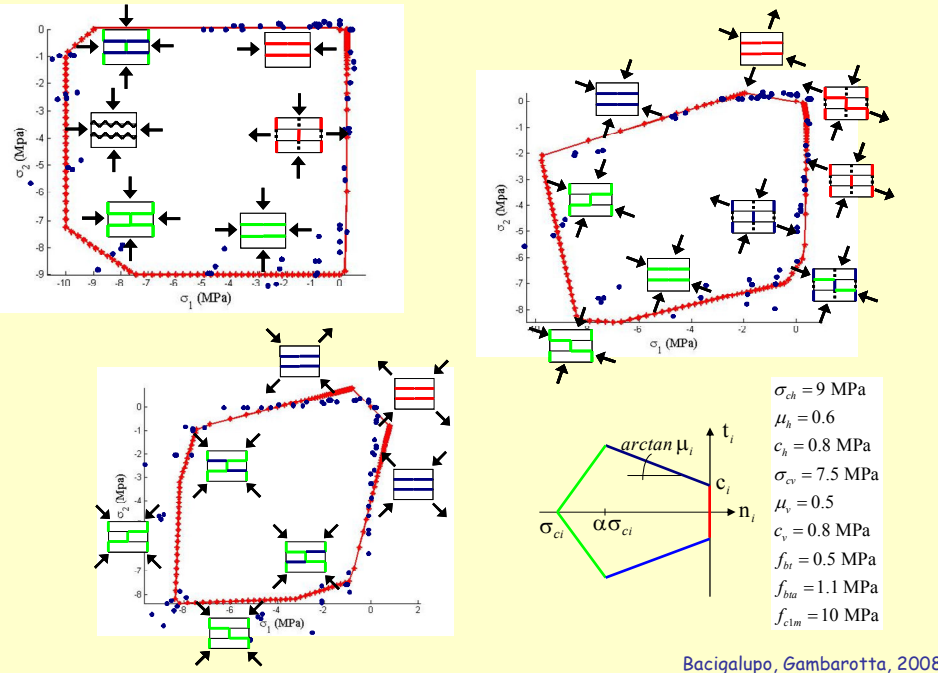


Strain localization

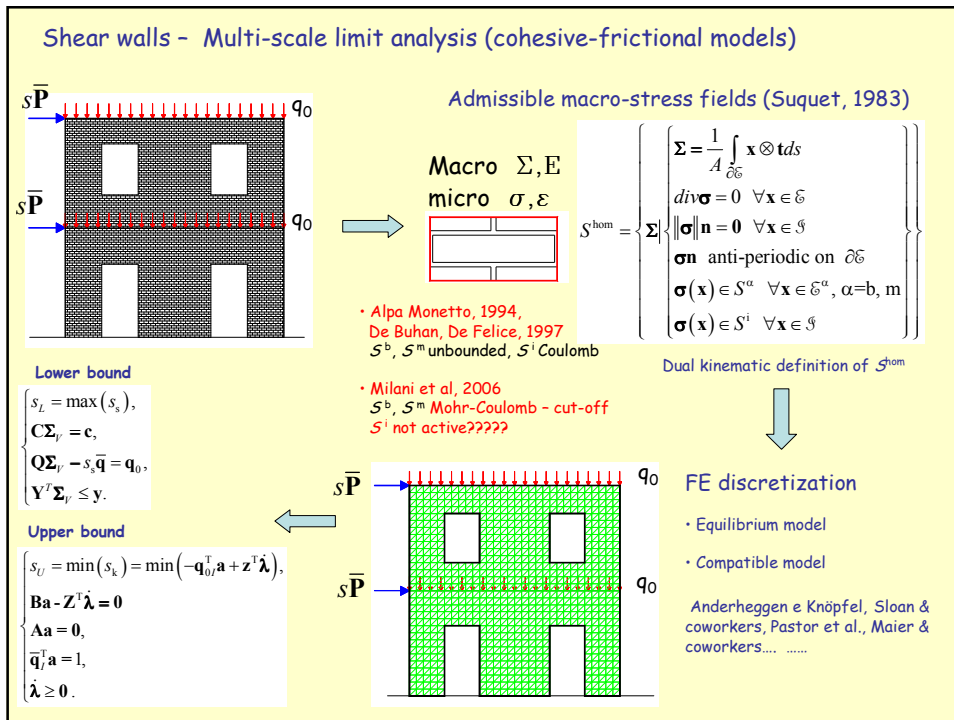
Limit strength domains by homogenization of the periodic unit cell



Limit strength domains: simulation of experimental results by Page



Shear walls - Multi-scale limit analysis (cohesive-frictional models)



Limit strength domains: simulation of experimental results by Page

12	11	10	9	8	7
15	14	13	12	11	10
18	17	16	15	14	13
21	20	19	18	17	16
24	23	22	21	20	19
27	26	25	24	23	22
30	29	28	27	26	25
33	32	31	30	29	28
36	35	34	33	32	31

Representative cell and sub-domains

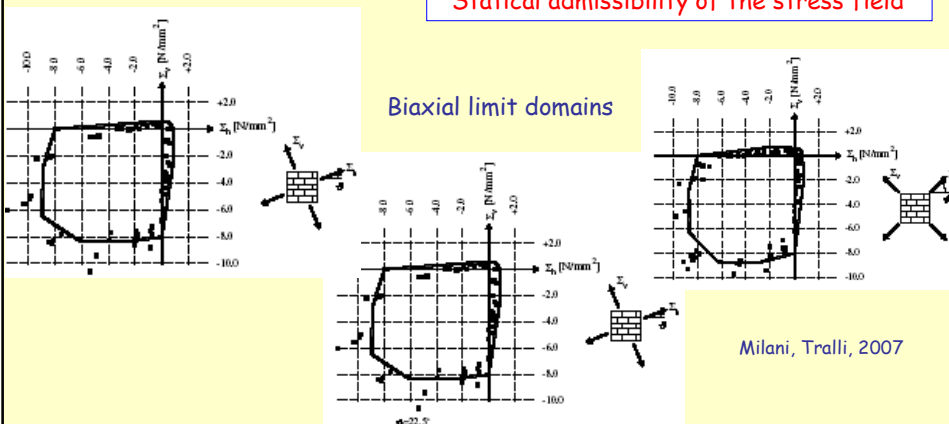
Polynomial distributions assumed for the stress components at each sub-domain

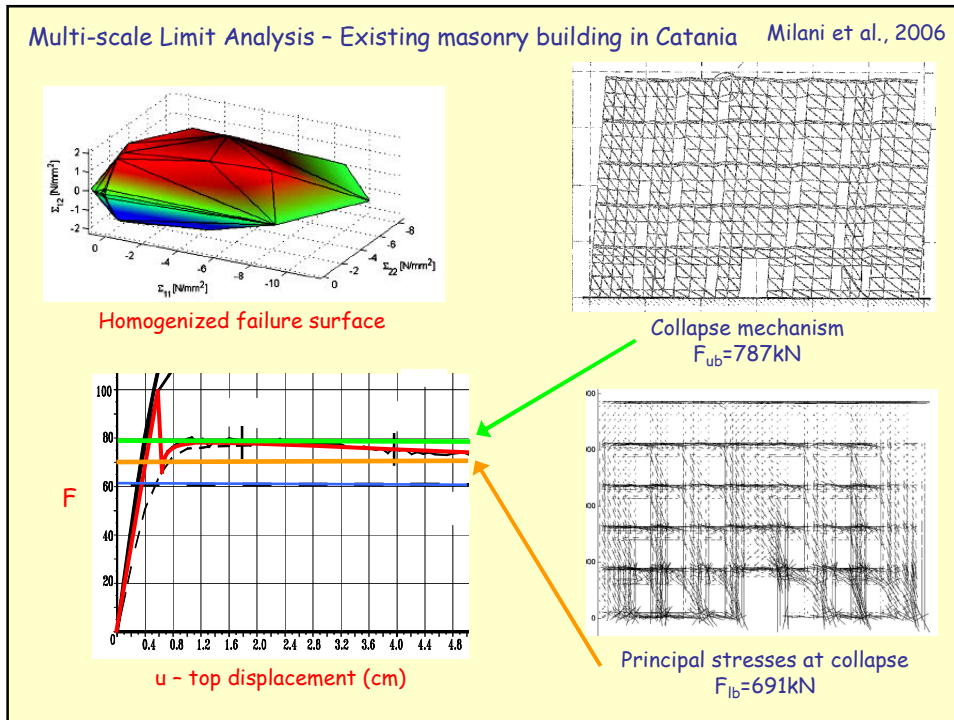
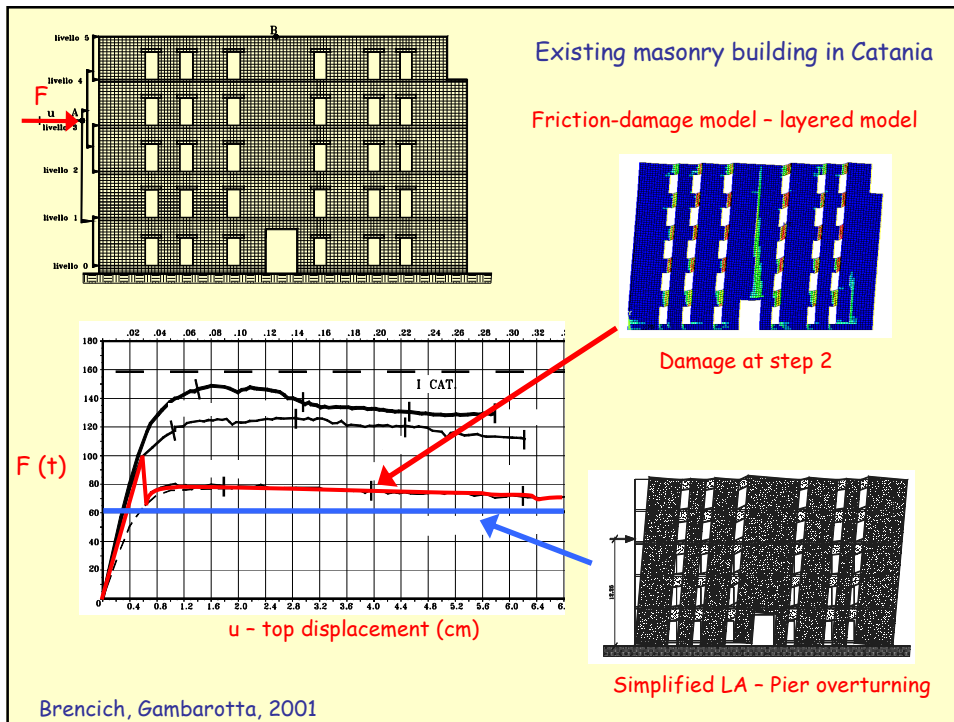
$$\boldsymbol{\sigma}_{ij}^{(k)} = \mathbf{z}(\mathbf{x})^T \mathbf{s}_{ij} \quad \mathbf{x} \in Y^k$$

Equilibrium and anti-periodicity of $\boldsymbol{\sigma} \mathbf{n}$ on $\partial \mathcal{E}$

Statical admissibility of the stress field

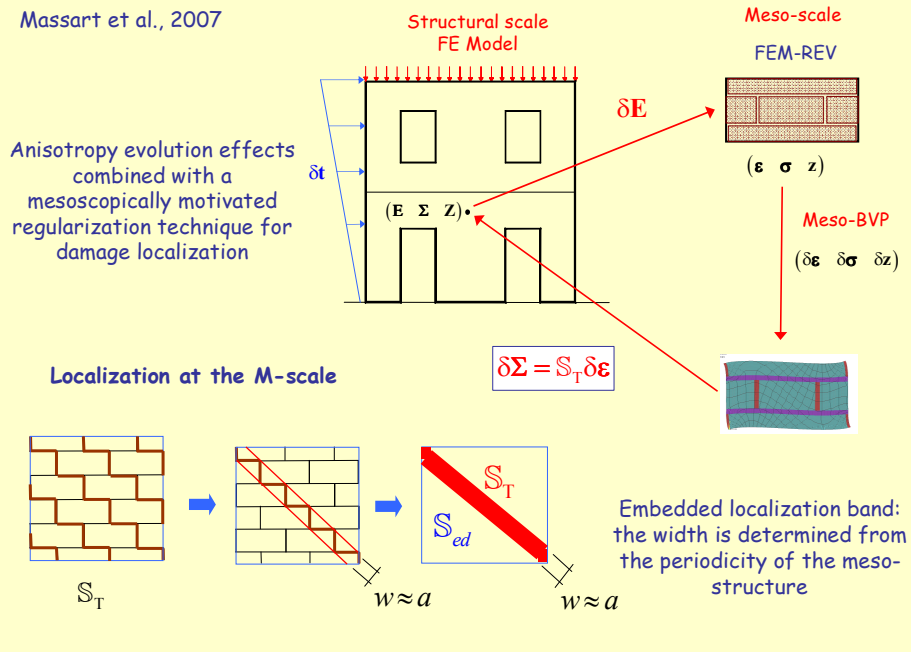
Biaxial limit domains





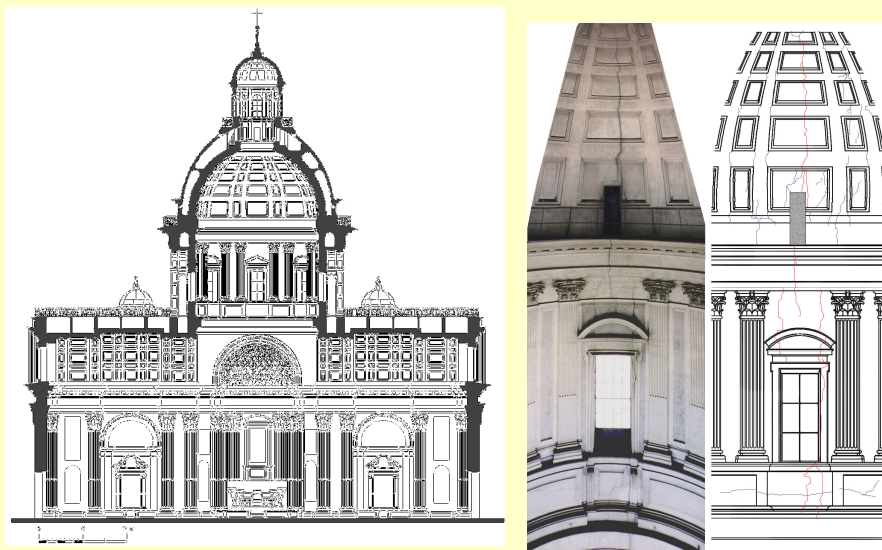
Multi-scale analysis of damaging shear walls

Massart et al., 2007

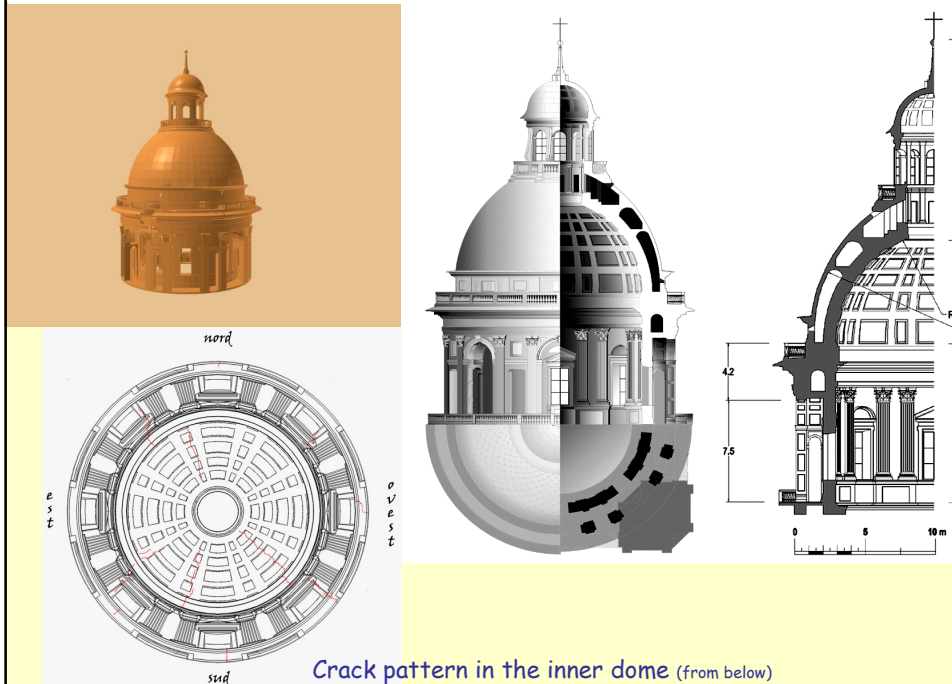


Masonry domes

Basilica di S. Maria di Carignano - Genova



Dome-drum interaction: Basilica di Carignano in Genova (G. Alessi, 1540-1600)



Crack pattern in the inner dome (from below)

Basilica di Carignano: Safe theorem
Statically admissible states

Hypotheses

- NTR material
- Infinite compressive strength
- No sliding failures admitted

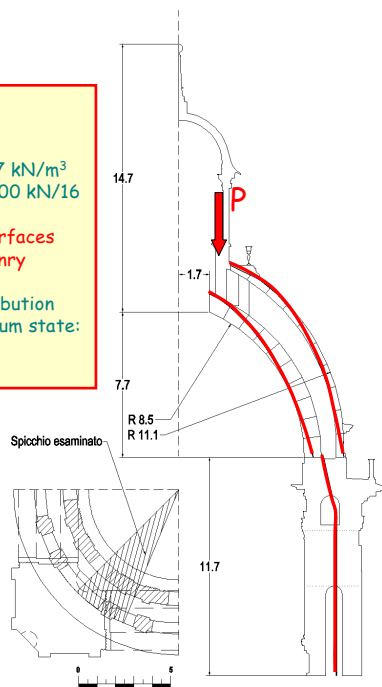
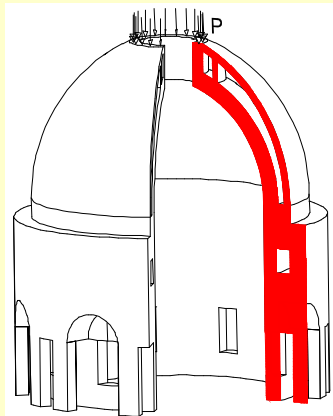
Equilibrium of a slice

Loads:

- masonry weight $\gamma=17 \text{ kN/m}^3$
- lantern weight $P=1200 \text{ kN/16}$

Search for thrust surfaces lying within the masonry

Lantern weight distribution for the safe equilibrium state:
85% inner shell
15% outer shell



Gambarotta et al., 2002

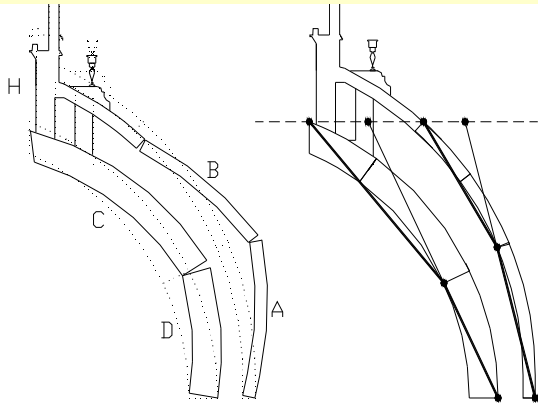
Upper Bound Theorem

If $\exists \dot{\mathbf{u}} \in \text{KinAdm}$ such that:

$$\dot{W} = \int_{\mathcal{B}^-} \mathbf{b} \cdot \dot{\mathbf{u}}^- dv + \int_{\mathcal{B}^+} \mathbf{b} \cdot \dot{\mathbf{u}}^+ dv = \dot{W}_a + \dot{W}_{res} \geq 0$$

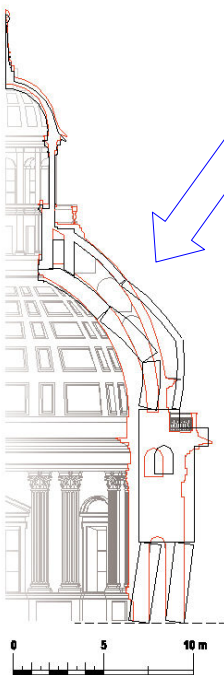
\Rightarrow The structure will collapse

- \mathbf{b} - unit volume weight
- \mathbf{u}^+ - upward velocity
- \mathbf{u}^- - downward velocity



1. Local mechanism
Inner and outer domes

$$\eta_1 = \frac{|\dot{W}_{res}|}{\dot{W}_a} \approx 2 > 1 \Rightarrow \dot{W} < 0$$



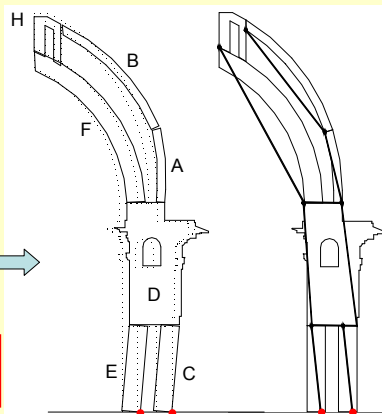
Global mechanism 1.

$$\eta_2 = \frac{|\dot{W}_{res}|}{\dot{W}_a} \approx 8.5$$

Global mechanism 2.

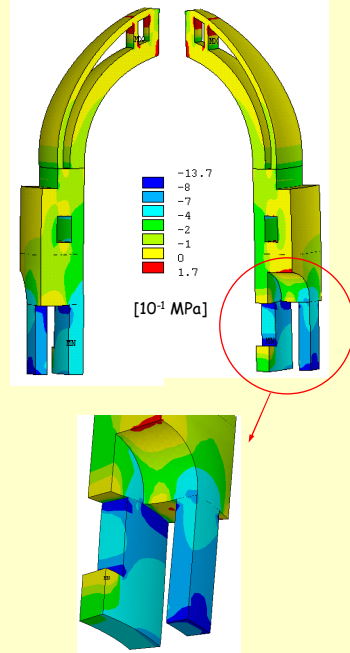
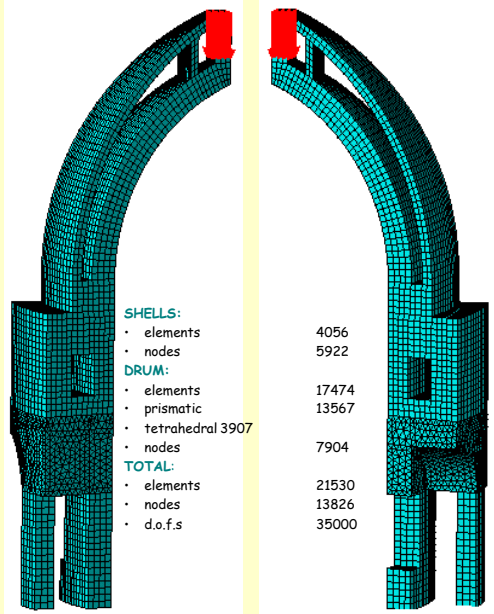
$$\eta_3 = \frac{|\dot{W}_{res}|}{\dot{W}_a} \approx 1.8 \div 7$$

Overall Mechanism domes-drum int.

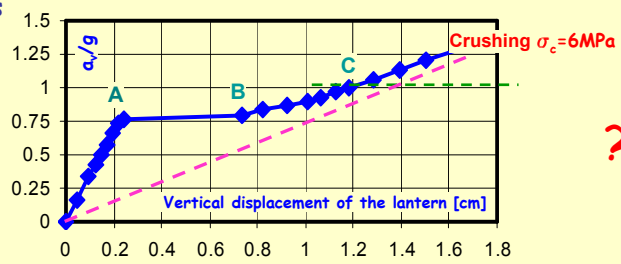


Influence of the compressive strength on the location of the centre of rotation of the drum slice

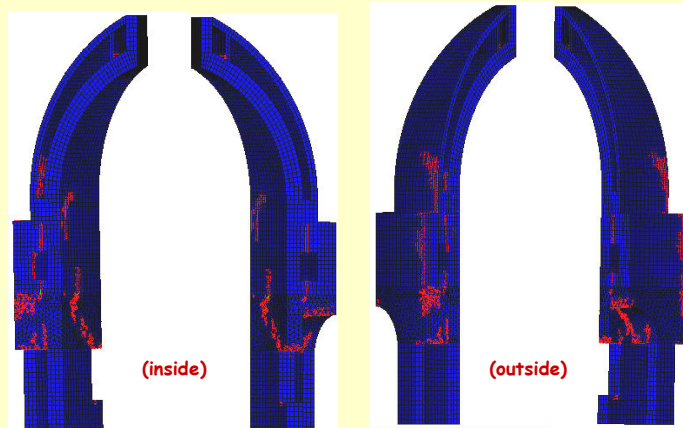
FE Model -1/8 slice



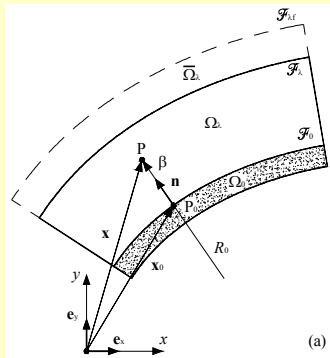
Incremental analysis



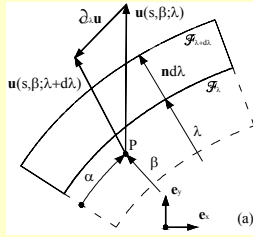
Crack pattern State A



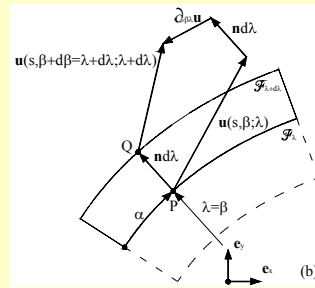
Influence of the construction sequence - structural growth



Reference domain



displacement increment due to growth



displacement increment accretion line

Stress field

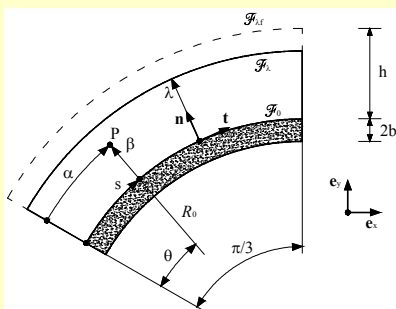
$$\mathbf{T}(s, \beta; \lambda_f) = \int_{\beta}^{\lambda_f} \mathbb{C} \text{sym} \nabla \mathbf{g}(s, \beta, \lambda) d\lambda$$

Strain field

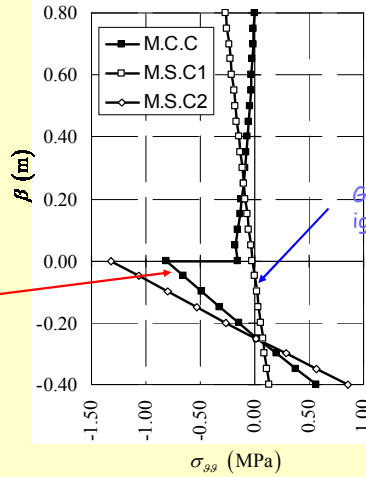
$$\mathbf{E}(s, \beta; \lambda) = \text{sym} \nabla \bar{\mathbf{u}}_0(s) + \int_0^{\beta} \text{sym} \nabla \bar{\mathbf{g}}(s, \beta = \lambda; \lambda) d\lambda + [\bar{\mathbf{g}}(s, \beta; \lambda = \beta) - \bar{\mathbf{g}}(s, \beta; \lambda = \beta)] \odot \mathbf{n} + \int_{\beta}^{\lambda} \text{sym} \nabla \bar{\mathbf{g}}(s, \beta; \lambda') d\lambda'$$

Influence of the construction sequence - structural growth

Example: Triumphal arch

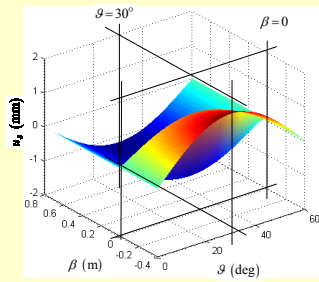


Normal stresses at crown section



Growth included

Growth ignored



Displacement field
Tangential component

Bacigalupo, Gambarotta, 2008

Acknowledgement

*Andrea, Bacigalupo, Antonio Brencich, Andrea Cavicchi, Christian Corradi,
Sergio Lagomarsino, Renata Morbiducci, Enrico Sterpi*

Thank You for Listening